 광주과학기술원
Gwangju Institute of Science and Technology

4. Magnetic Resonance Imaging(1)

Lectures 17, 18

Medical Imaging Systems


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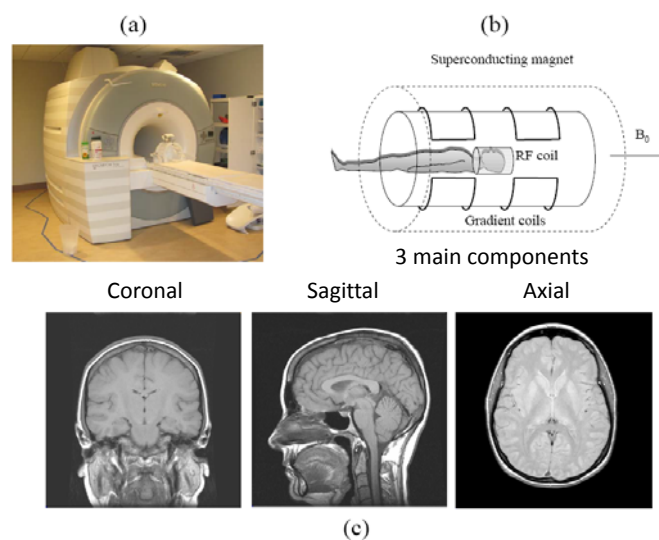
Contents

1. Introduction
2. Effects of a strong magnetic field on protons in the body
3. Effects of radiofrequency pulse on magnetization
4. Faraday induction: the basis of MR signal detection
5. T1 and T2 relaxation times
6. Signals from lipid
7. The free induction decay

Introduction

- The first MRI image: Paul Lauterbur in 1973 → won Nobel prize for medicine in 2003 with Peter Mansfield
- Advantages
 - Non-ionizing radiation
 - Images can be acquired in any 2-3D plane
 - Excellent soft-tissue contrast
 - Good spatial resolution (< 1mm)
 - Can provide information other than anatomy such as blood flow and water diffusion
- Disadvantages
 - Acquisition is much slower than CT/Ultrasound, comparable to PET
 - Need to exclude patients who has metallic implants
 - Much expensive than CT or ultrasound units

Introduction



Video Lectures

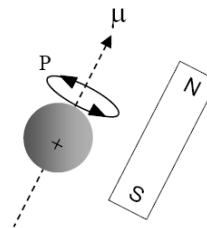
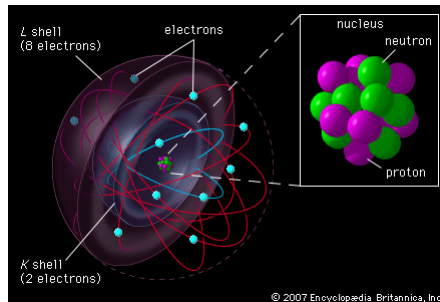
- <http://youtu.be/7aRKAXD4dAg>
- <http://youtu.be/jUKdVBpCLHM>
- <http://youtu.be/GjLvu1hOAAA>

Magnetic field effects on protons

- A typical value of magnetic field, B_0 , is 3T (30,000Gauss)
- Examples of magnetic field value
 - Earth: $30\mu\text{T}$ at equator, $70\mu\text{T}$ at poles
 - Refrigerator: 10mT
 - Clinical MRI: 1.5~3T
 - Research NMR: 4.7~11T +
- Quantum mechanical description
- Classical description

Quantum Mechanical Description

- A nuclei that has odd number of protons, neutrons, or both spins around an internal axis of rotation with a given value of **angular momentum (P)** or called 'spin'.
- Hydrogen atom is one of them (a single proton)
- This angular momentum produces a **magnetic moment (μ)**
- Now, this nuclei acts like a small bar magnet



Quantum Mechanical Description

- The magnitude of **angular momentum** of the proton is quantized (i.e. it can only have discrete values and spin quantum number, I , is $\frac{1}{2}$ for a hydrogen).

$$|\vec{P}| = \frac{h}{2\pi} [I(I + 1)]^{1/2}$$

- Therefore, $|\vec{P}| = \frac{h}{2\pi} \frac{\sqrt{3}}{2}$ for a hydrogen
- The proton's **magnetic moment** is proportional to the magnitude of angular momentum

$$|\vec{\mu}| = \gamma |\vec{P}|$$

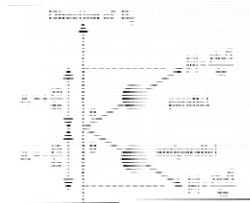
where γ is gyromagnetic ratio (42.6 MHz/Tesla for protons)

- Therefore, the magnetic moment also has a single, fixed value

$$|\vec{\mu}| = \frac{\gamma h}{2\pi} [I(I + 1)]^{1/2}$$

Quantum Mechanical Description

- Therefore, $|\vec{\mu}| = \frac{\gamma h \sqrt{3}}{4\pi}$ for a hydrogen
- With a strong magnetic field, μ will have only the values given by $\mu_z = \gamma P_z = \frac{\gamma h}{2\pi} m_l$ where m_l is the nuclear magnetic quantum number, and can take values, $l, l-1, \dots, -l$, For hydrogen, l is $\frac{1}{2}$ and $-\frac{1}{2}$
- Therefore, $\mu_z = \pm \frac{\gamma h}{4\pi}$ and this makes a proton rotates with an angle of 54.7° ($\cos\theta = \frac{\frac{\gamma h}{4\pi}}{\frac{\gamma h \sqrt{3}}{4\pi}} = \frac{1}{\sqrt{3}}$, $\theta = \cos^{-1} \frac{1}{\sqrt{3}} = 54.74^\circ$)



Quantum Mechanical Description

- The energy of i th spin state (E_i) is directly proportional to the value of m_l (nuclear magnetic quantum number) and magnetic field strength B_0

$$E_i = -\mu_z B_0 = -m_l \frac{\gamma h B_0}{2\pi} = \mp \frac{\gamma h B_0}{4\pi} \text{ (for hydrogen)}$$

- The energy difference between anti-parallel and parallel (ΔE) is

$$\Delta E = \left| E_{-\frac{1}{2}} - E_{\frac{1}{2}} \right| = \left| \left[\left(\frac{1}{2} \right) \left(\frac{\gamma h B_0}{2\pi} \right) \right] - \left[- \left(\frac{1}{2} \right) \left(\frac{\gamma h B_0}{2\pi} \right) \right] \right| = \frac{\gamma h B_0}{2\pi}$$

where h is Plank's constant (6.63×10^{-34} J s)

Quantum Mechanical Description

- The relative number of nuclei in each configuration can be obtained by using a Boltzmann equation.

$$\frac{N_{\text{antiparallel}}}{N_{\text{parallel}}} = \exp\left(-\frac{\Delta E}{kT}\right) = \exp\left(-\frac{\gamma h B_0}{2\pi kT}\right)$$

where k is the Boltzmann coefficient (1.38×10^{-23} J/K)

- A first-order approximation of $e^{-x} \approx 1 - x$
- Therefore, $\frac{N_{\text{antiparallel}}}{N_{\text{parallel}}} = 1 - \frac{\gamma h B_0}{2\pi kT}$ and

$$N_{\text{parallel}} - N_{\text{antiparallel}} = N_{\text{total}} \frac{\gamma h B_0}{4\pi kT}$$

where N_{total} is the total number of protons in the body

- In 1.5T and 1 million protons, only 5 protons difference between parallel and antiparallel orientation

Quantum Mechanical Description

- We should note that if $I=0$, then there is no angular momentum and thus no magnetic moment \rightarrow no MR signal

TABLE 4.1. Properties of Nuclei Found at High Abundance in the Body

Nucleus	Atomic Number	Atomic Mass	I	$\gamma/2\pi$ (MHz/T)	MRI Signal
Proton	1	1	1/2	42.58	Yes
Phosphorus	15	31	1/2	17.24	Yes
Carbon	6	12	0	—	No
Oxygen	8	16	0	—	No
Sodium	11	23	3/2	11.26	Yes

- That's why we use ^{13}C and ^{17}O for MRI

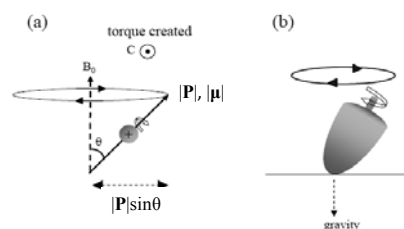
Classical Description

- We described nuclear magnetism using the quantum mechanical model, but it becomes harder to analyze the complicated MRI pulse sequences → use classical mechanics
- When protons are placed in the external magnetic field (B_0),
 - the B_0 field attempts to align the proton magnetic moment with itself
 - This action creates a torque, C

$$\vec{C} = \vec{\mu} \times \vec{B}_0 = i_N |\mu| |B_0| \sin\theta$$

where i_N is a unit vector normal to both $\vec{\mu}$ and \vec{B}_0

Classical Description

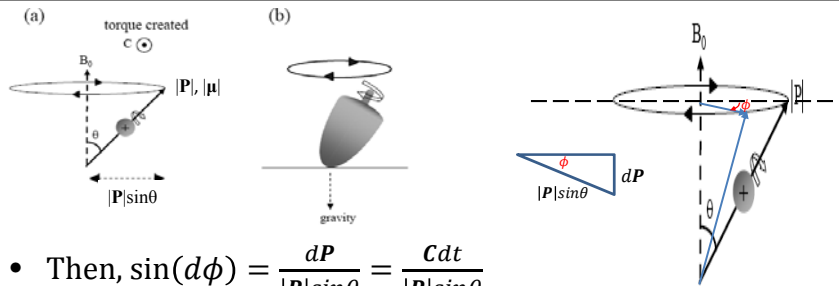


- In order to calculate the proton precession frequency, consider that the torque is defined as follows;

$$\vec{C} = \frac{d\vec{P}}{dt} = \vec{\mu} \times \vec{B}_0$$

- The magnitude of angular momentum perpendicular to B_0 is $|P| \sin\theta$
- In a short time dt , the magnetic moment precesses through an angle $d\phi$ producing a change dP in the angular momentum

Classical Description

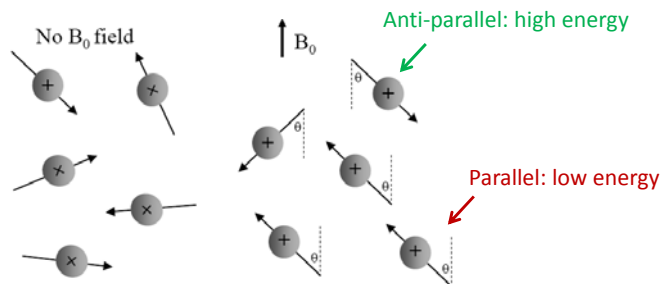


- Then, $\sin(d\phi) = \frac{dP}{|P|\sin\theta} = \frac{Cdt}{|P|\sin\theta}$
- If $d\phi$ is small, then $\sin(d\phi) \sim d\phi$
- The angular precession frequency ω is given by $d\phi/dt$
- Therefore,

$$\omega = \frac{d\phi}{dt} = \frac{Cdt}{dt|P|\sin\theta} = \frac{\mu \times B_0}{|P|\sin\theta} = \frac{\gamma P \times B_0}{|P|\sin\theta} = \frac{\gamma |P| |B_0| \sin\theta}{|P|\sin\theta} = \gamma B_0$$
 where B_0 represent the magnitude of B_0

B_0 Field Effect

- For the normal state, the orientation of magnetic moment of protons are random \rightarrow net magnetic moment is zero
- If these protons are placed in the external magnetic field (B_0), they start to “precess” around the axis of B_0 with an angle of $54.7^\circ \rightarrow$ Larmor frequency (ω) = γB_0

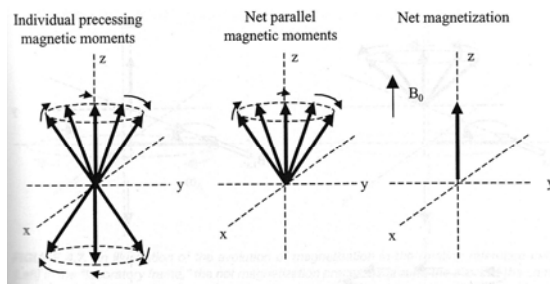


B₀ Field Effect

- Net magnetization of a sample is defined as M₀:

$$M_0 = \sum_{n=1}^{N_{total}} \mu_{z,n} = \frac{\gamma h}{2\pi} (N_{parallel} - N_{antiparallel}) = \frac{\gamma^2 h^2 B_0 N_{total}}{2\pi^2 kT}$$

- To have a larger M₀,
→ higher γ, higher B₀, lower T



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Proton energy levels

No B₀ field

no magnetic field

B₀ present

anti-parallel
parallel

↑ B₀

Anti-parallel: high energy

Parallel: low energy

$E = +\frac{\gamma h B_0}{4\pi}$

$E = -\frac{\gamma h B_0}{4\pi}$

ΔE

B₀ present

Zeeman effect

Energy Levels

Spectra

No Magnetic Field Magnetic Field

The number of split levels in the magnetic field is $2 * L + 1$ (L: spin quantum number)

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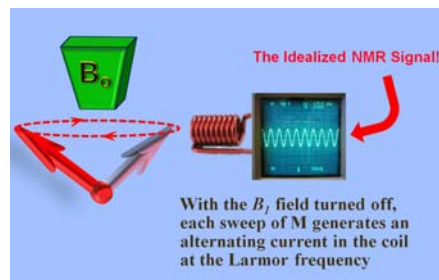
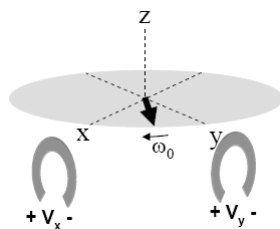
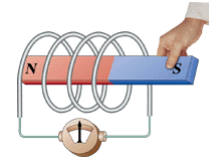
Faraday induction

- Faraday's law of induction: the voltage (V) is induced in each of the loop with a value proportional to the time rate of change of the magnetic flux $d\phi$:

$$V = -n \frac{d\phi}{dt} \quad n: \text{number of coils}$$

$$V_y \propto M_0 \omega_0 \sin \omega_0 t,$$

$$V_x \propto -M_0 \omega_0 \cos \omega_0 t$$



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RF pulse effects on magnetization

- From the previous slide, we can see that only oscillating magnetic field can be observed
- However, M_z is not oscillating
- Therefore, we need to rotate precessing protons from z axis into xy plane
- Solution:** apply a second magnetic field at a 90° angle to the z axis \rightarrow creates a second torque \rightarrow rotates the magnetization toward the xy plane
- This second magnetic field can be generated from a RF pulse

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RF pulse effects on magnetization

- Energy levels of protons in magnetic field are analogous to
 - Energy levels in semiconductors
 - Vibrational and rotational energy levels in infrared spectroscopy
- To obtain an MR signal, **energy must be supplied with a specific value ΔE to stimulate transitions between energy levels**
- De Broglie's relationship of frequency: $\Delta E = hf$:

$$hf = \Delta E = \frac{\gamma h B_0}{2\pi}, \quad f = \frac{\gamma B_0}{2\pi}, \quad \omega = \gamma B_0 \text{ (Larmor frequency)}$$

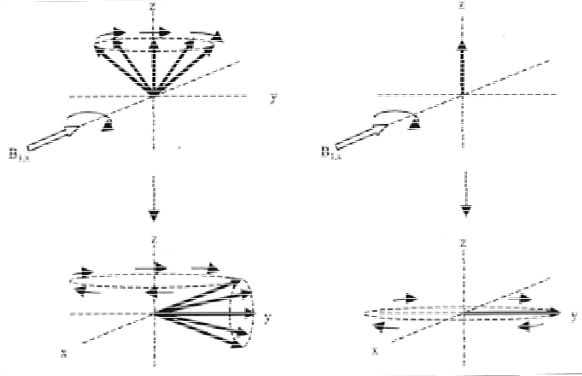
→ *Resonance frequency*

RF pulse effects on magnetization

- Transverse magnetization creation
 - In MRI, energy is applied as a short RF pulse 90° to B_0 (creates B_1 field)
 - B_1 field produces a torque which causes the net magnetization to rotated towards the xy -plane
 - The tip angle (α) can be obtained by

$$\alpha = \gamma B_1 \tau_{B1}$$

RF pulse effects on magnetization



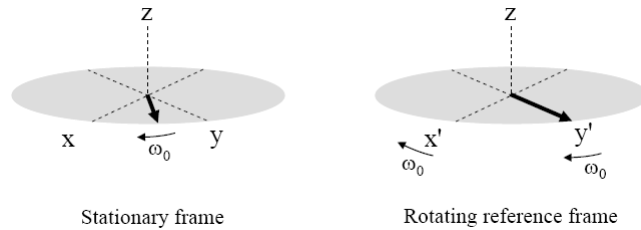
- (A) Application of B_1 field along the x axis rotates the individual proton magnetic moments around the x axis toward y axis, (B) After applying the B_1 field for a certain time duration, the “cone” of magnetic moments has been rotated by 90° . The magnetic moments continue to precess around the B_0 axis, (C) and (D) show the vector model representation of (A) and (B)

MR signal intensity

- From $V_y \propto M_0 \omega_0 \sin \omega_0 t$, $V_x \propto M_0 \omega_0 \cos \omega_0 t$, the size of MR signal is depending on
 - The number of protons; $M_0 = \frac{\gamma^2 \hbar^2 B_0 N_{total}}{2\pi^2 kT}$
 - The value of B_0 ; $M_0 = \frac{\gamma^2 \hbar^2 B_0 N_{total}}{2\pi^2 kT}$
 - The precession frequency, ω_0 (\rightarrow value of B_0)
- Overall, the image SNR $\propto B_0^{3/2}$

The rotating reference frame

- To visualize the proton precession, the concept of a 'rotating reference frame (x', y', z')' is useful

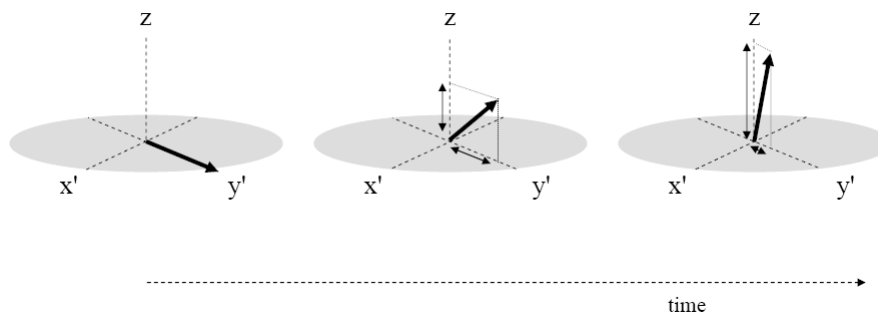


T1 relaxation times (spin-lattice)

$$M_z(t) = M_0 \cos\alpha + (M_0 - M_0 \cos\alpha) \left(1 - e^{-\frac{t}{T_1}}\right)$$

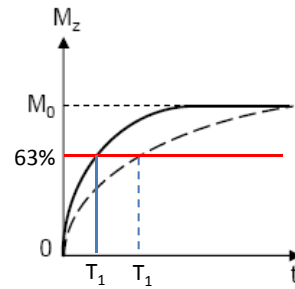
- After a 90° pulse,

$$M_z(t) = M_0 \left(1 - e^{-\frac{t}{T_1}}\right)$$



T1 relaxation times

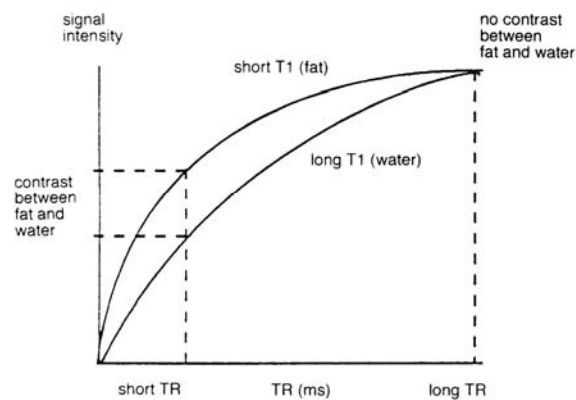
- The recovery of M_z magnetization after a 90° pulse



- When $t=5T_1$, $M_z \sim 99\%M_0$, $t=3T_1$, $M_z \sim 95\%M_0$

T1 Weighting

- TR controls the amount of T1 weighting
- For T1 weighting the TR must be short

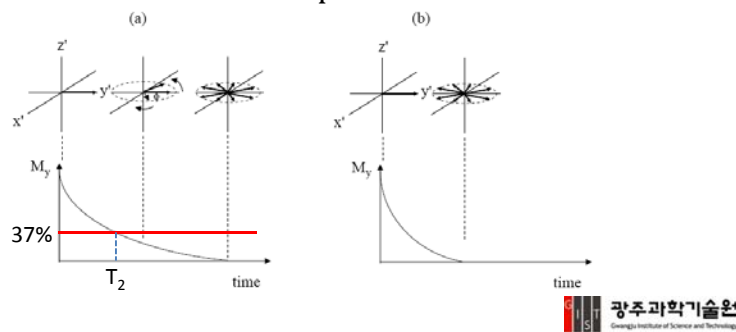


T2 relaxation times (spin-spin)

- If RF pulse is applied along the x-axis, M_y at time t after RF pulse is :

$$M_y(t) = M_0 \cdot \sin\alpha \cdot e^{-\frac{t}{T_2}}$$

In practice, **molecular dynamics (interactions with neighboring nuclei) causes a small spread in the precessional frequencies** although all protons are assumed to precess at the same frequency



T2 relaxation times

- Additional factor that leads M_x and M_y back to zero
 - **Spatial inhomogeneity's of B_0 field**
 - Impossible to design a magnet which produces a perfect uniform B_0
 - Local variations in magnetic field due to the different magnetic susceptibilities of different parts of the body (eg. air/tissue, bone/tissue boundaries)

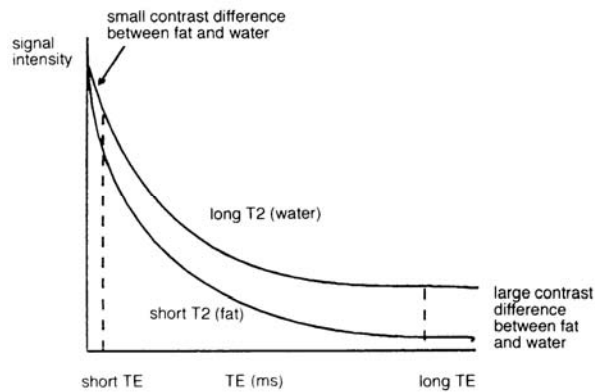
– T_2^+ : the effects of local B_0 field inhomogeneity

$$\frac{1}{T_2^*} = \frac{1}{T_2^+} + \frac{1}{T_2}$$

– In high resolution NMR, $T_2^* \cong T_2$ since T_2^+ is very small (sample is small and M_z is spatially homogeneous)

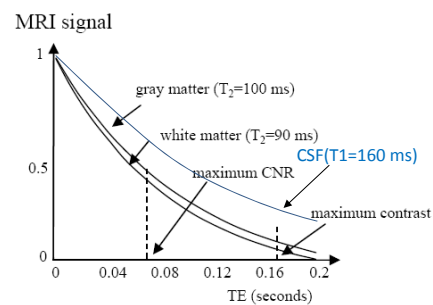
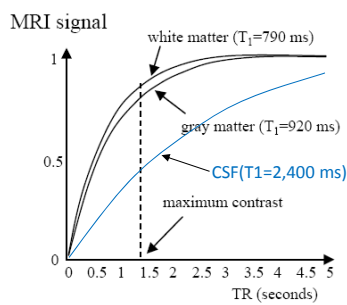
T2 Weighting

- TE controls the amount of T2 weighting
- For T2 weighting the TE must be long



Brain T_1 and T_2

- MRI signal is shown with 1.5T

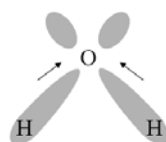


Video Tutorials

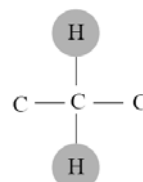
- 90 degree pulse and T1 relaxation
- <http://youtu.be/P0YuwCphcRU>
- T2 dephasing
- <http://youtu.be/lKp67lqQjH4>
- T2 dephasing topview
- http://youtu.be/_7oZMA0OuK4

Signals from lipid

- The majority of protons in lipid are in the form of $-CH_2-$ groups in long-chain fatty acids
- The electron density distribution (shaded area) surrounding protons in water and lipid are shown in figures below
- Proton in water is less shielded due to strong electronegativity of oxygen than carbon in lipid



water



Fatty acid

Signals from lipid

- An electron is a negatively charged particle, so produces a small magnetic field opposite in polarity to the B_0

$$B_{\text{eff}} = B_0(1 - \sigma)$$

where σ is a shielding constant

- Therefore, the resonant frequency of the proton is

$$\omega = \gamma B_{\text{eff}} = \gamma B_0 (1 - \sigma)$$

- σ for water is less than that for lipid, ($\sigma_{\text{water}} < \sigma_{\text{lipid}}$) therefore, water has a higher resonant frequency ($B_{\text{eff, water}} > B_{\text{eff, lipid}}$) (at 3T the difference is about 2500 rad/s or 400Hz)

The free induction decay

- The signal precesses freely after the RF pulse has been turned off
- This M_x and M_y shows beat patterns which come from the two different resonant frequencies of lipid and water

